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Ionic Conductivity in Lithium Hexaoxometallate

Solid Solutions

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20. ABSTRACT (Continue on reverse side II necessary and identity by block number)

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are 4.3 x 10⁻⁴ (Ωcm)⁻¹, 3.0 x 10⁻⁴ (Ωcm)⁻¹ and 4.0 x 10⁻⁴ (Ωcm)⁻¹,

respectively.

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IONIC CONDUCTIVITY IN LITHIUM HEXAOXOMETALLATE SOLID SOLUTIONS

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ABSTRACT

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Lithium ion conductivity in solid solutions formed between Li₇Ta₀₆, Li₇M^VO₆ (M^V = Nb, Bi) and Li₈ZrO₆ has been measured as a function of temperature and composition using complex impedance method. The samples prepared were Li₇Ta_{1-x}Nb_xO₆, Li₇Ta_{1-x}Bi_xO₆ and Li_{7+x}Ta_{1-x}Zr_xO₆. At 200°C, the conductivities of Li₇Ta_{0.7}Nb_{0.3}O₆, Li₇Ta_{0.6}Bi_{0.4}O₆ and Li_{7.4}Ta_{0.6}Zr_{0.4}O₆ are 4.3 x 10^{-4} (Ω cm)⁻¹, 3.0 x 10^{-4} (Ω cm)⁻¹ and 4.0 x 10^{-4} (Ω cm)⁻¹ respectively.

INTRODUCTION

There is considerable interest in developing solid lithium ion conductors for utilization in high energy density battery systems. This has led to a search for new solid electrolytes exhibiting high lithium ion conductivity and has stimulated interest in developing a fundamental understanding of ionic transport in solids (1-3).

The compounds of lithium hexaoxometallates, formulated as $\text{Li}_{1}\text{MO}_{6}$; n = 6,7 or 8, M = IV, V or VI group element, have a pseudo two dimensional structure and high lithium ion conductivity⁽⁴⁾. The structure of $\text{Li}_{1}\text{MO}_{6}$ is characterized by octahedral sheets of CdI_{2} -type, between which 6 Li⁺ are inserted in a tetrahedral environment as

a+b+c = 3

: vacancy

Among the lithium hexaoxometallates, Li₇TaO₆ (i.e. Li₆(LiTaD)O₆) is the best lithium conductor with σ_{200} = 5 x 10⁻⁴ (Ω cm)⁻¹ and σ_{RT} = 4.3 x 10⁻⁸ (Ω cm)⁻¹(4).

It has been shown that the activation energy for Li⁺ diffusion in the tetrahedral sites between the layers is significantly larger than that for the octahedral lithium ion in the layers (5). Thus by altering the lithium content and/or the vacant sites in the layers higher ionic conductivities might be achieved. We have examined the effect of increased Li⁺ content by partial substitution of tantalum ions with Zr^{4+} to yield $Li_6(Li_{1+x}Ta_{1-x}Zr_x \ \square_{1-x})0_6$. We have also attempted to look at the effect of decreased Li⁺ content (or increased vacancies) in the layers by studying the conductivities of samples with substituted divalent cations (e.g. $Li_6(Li_{1-2x}Mg_xTa\ \square_{1+x})0_6$), however our results are inconclusive on this system due to difficulties in obtaining single phase compounds.

Alternatively we hoped to obtain higher lithium conductivities in this system by partial substitution of the tantalum ions with other group V metal ions of larger ionic radii in order to optimize the channel size of the framework structure for lithium diffusion. In this paper we report the results of the lithium conductivity in: $\text{Li}_7\text{Ta}_{1-x}\text{Nb}_x\text{O}_6$, $\text{Li}_7\text{Ta}_{1-x}\text{Bi}_x\text{O}_6$, $\text{Li}_7+_x\text{Ta}_{1-x}\text{Zr}_x\text{O}_6$.

EXPERIMENTAL

Li₂O was obtained by thermal decomposition of anhydrous Li₂O₂ in vacuum at 450°C for 6 hours. Other starting materials were reagent grade Nb₂O₅, Ta₂O₅, Bi₂O₃ and ZrO₂. Mixtures of appropriate composition were throughly mixed using agate mortar in a He dry box. For example, Li₇Ta_{1-x}Nb_xO₆ and Li₇Ta_{1-x}Bi_xO₆ were prepared from Li₂O, Ta₂O₅ and Nb₂O₅, and Li₂O, Ta₂O₅ and Bi₂O₃, respectively, according to the following equation

$$7\text{Li}_2\text{O} + (1-x)\text{Ta}_2\text{O}_5 + x\text{Nb}_2\text{O}_5 + 2\text{Li}_7\text{Ta}_{1-x}\text{Nb}_x\text{O}_6$$

 $7\text{Li}_2\text{O} + (1-x)\text{Ta}_2\text{O}_5 + x\text{Bi}_2\text{O}_3 + x\text{O}_2 + 2\text{Li}_7\text{Ta}_{1-x}\text{Bi}_x\text{O}_6^{(6)}$.

The mixtures were pressed into cylindrical pellets 3/8" diameter and about 0.2" thick at 15 kpsi. The pressed pellets were transferred to high purity alumina crucibles, were embedded in excess Li₂O powder in order to prevent Li₂O losses, and were heated at 650°C for 18 hours in air. After cooling, the pellets were crushed and examined by X-ray powder diffraction using a Norelco diffractometer with Ni filtered copper radiation. Lithium content was determined by atomic absorption spectrophotometric methods.

Pellet samples for ionic conductivity measurements were prepared by pressing to 1/4" diameter and about 0.15" thickness at 30 kpsi followed by sintering at 950°C for 18 hours in air and quenching in air. Again, Li₂0 loss was minimized by covering the pellets with Li₂0 powder during the sintering process. The X-ray diffraction pattern and lithium content of the sintered samples were checked to confirm the identity and composition of the phases present. Weight loss was not observed after sintering. The density of the sintered pellets were about 85% of theoretical value. Both surfaces of the pellets were polished using silicon carbide (#400) paper and sputtered with about lµm of gold followed by a coating of silver paint (Engelhard #16).

The device for the conductivity measurement is shown in Fig. 1. Platinum contact leads were made by spot welding pieces of platinum foil of $10 \times 10 \text{ mm}^2$ to platinum wires. Two discs of α -alumina of 10 mm diameter and 3 mm thickness were used for insulating the materials. Contact between the sample electrode (sputtered gold and painted silver) and the platinum lead was made

Fig. 1 Device for conductivity measurement

 sample Pellet
 platinum lead
 draumina disc
 stainless steel disc 5. screw

Fig. 2

Circuit diagram for AC conductivity measurement

A. tuned amplifier PS. phase shifter M. mixer I. integrator

R_x. sample R_{std}. standard resistance by tightening the screw (Fig. 1). The AC conductivity measurements were made using ionically blocking electrodes. The schematic electrical circuit diagram of the AC conductivity measurement is shown in Fig. 2. Lock-in amplifier (PAR Model 128A) was used as a phase sensitive detector and a Hewlett-Packard Model 200 CDR was used as a signal generator. The frequency range used was between 5 HZ and 100 KHZ. The measurement temperature was from room temperature to 250°C in air.

RESULTS AND DISCUSSIONS

Li7Ta06

The ionic conductivity of $\sigma_{200} \cong 5 \times 10^{-4}~(\Omega~cm)^{-1}$ and activation energy of 0.66 eV of Li7TaO₆ have been measured by C. Delmas et al⁽⁴⁾. We have also measured the conductivity of sintered pellets of Li7TaO₆. Fig. 3 shows the conductivity, σ versus $^1/_T$ of Li7TaO₆. Different activation energies corresponding to a lower and higher temperature regions are observed. In the higher temperature region, the deviation of ionic conductivity between different pellets is small (which is also a measure of the reliability of the data) and the activation energy is 0.67 eV. In the lower temperature region, the ionic conductivity of different pellets differ more significantly but the activation energies are the same (0.46 eV). The change in the slope of the σ vs $^1/_T$ plot of Li7TaO₆ probably corresponds to a transition from extrinsic to intrinsic conductivity regions; then the variations in the ionic conductivities of different samples in the lower temperature region can be explained by differences in defect concentrations or particle packing from sample to sample.

Fig. 3

The temperature dependence of the conductivity of Li₇TaO₆

The X-ray diffraction patterns of the solid solutions prepared were identical to that of Li7Ta06 except for the shifts in 20. The lattice parameters calculated for a hexagonal unit cell increased with increasing values of x for Li7Ta1-xNbx06 and for Li7Ta1-xBix06; a = 5.42Å, c = 15.17Å for Li7Ta0.6Nb0.406 and a = 5.44Å c = 15.23Å for Li7Ta0.6Bi0.406. The lattice parameters of Li7Ta06, Li7Nb06 and Li7Bi06 are a = 5.39Å; c = 15.11Å; a = 5.40Å; c = 15.12Å; and a =5.50Å, c = 15.45Å, respectively(7). It is clear that the lattice parameters of Li7Ta0.6Nb0.406 are larger than those of either of the component compounds, Li7Ta06 and Li7Nb06, while the lattice parameters of Li7Ta0.6Bi0.406 are intermediate between those of Li7Ta06 and Li7Bi06. The variation of lattice parameters with x in Li7Ta1-xNbx06 and Li7Ta1-xBix06 is shown in Fig. 4. The bismuth substituted phases appear to obey Vagard's law while the Li7Ta1-xNbx06 system does not. This suggests that there might be a new phase close to the Li7Ta0.6Nb0.406 composition.

The ionic conductivities of $\text{Li}_{7}\text{Ta}_{1-x}\text{Nb}_{x}0_{6}$ are shown in Fig. 5. The ionic conductivities increase from 1.1 x 10^{-4} (Ω cm)⁻¹ for x = 0.0 to 4.3 x 10^{-4} (Ω cm)⁻¹ for x = 0.3 at 200°C and decrease as x is increased further (Fig. 5). The activation energy of conduction is 0.63 eV for $\text{Li}_{7}\text{Ta}_{0.7}\text{Nb}_{0.3}0_{6}$, somewhat smaller than the 0.67 eV observed for $\text{Li}_{7}\text{Ta}0_{6}$.

The ionic conductivities of $\text{Li}_7\text{Ta}_{1-x}\text{Bi}_x\text{O}_6$ phases are shown in Fig. 6. The ionic conductivities increase to 3.0 x 10^{-4} (Ωcm)⁻¹ at 200°C at x = 0.4. The activation energy for conduction is 0.60 eV in Li $_7\text{Ta}_{0.7}\text{Bi}_{0.3}\text{O}_6$.

Fig. 4
Lattice parameters of Li₇Ta_{1-x}Nb_xO₆(0) and Li₇Ta_{1-x}Bi_xO₆ (0)

Fig. 5 The temperature dependence of the conductivity of $\text{Li}_7\text{Ta}_{1-x}\text{Nb}_x\text{O}_6$ phases

Fig. 6 The temperature dependence of the conductivity of $\text{Li}_7\text{Ta}_{1-x}\text{Bi}_x\text{O}_6$ phases

In the Li7Ta_{l-x}Nb_xO₆ and Li7Ta_{l-x}Bi_xO₆ phases the observed increase in ionic conductivity and decrease in activation energy at x = 0.3 and 0.4 respectively compared to Li7TaO₆ is clearly due to optimizing the channel size for lithium diffusion at these compositions.

Li7+xTal-xZrx06

The powder X-ray diffraction pattern of $\text{Li}_{7+x}\text{Ta}_{1-x}\text{Zr}_x0_6$ is similar to that of Li₇Ta_{0.6}. The lattice parameters increased with increasing values of x; a = 5.44Å and c = 15.25Å were found for Li_{7.4}Ta_{0.6}Zr_{0.4}0₆.

The ionic conductivities of $\text{Li}_{7+x}\text{Ta}_{1-x}\text{Zr}_x0_6$ are shown in Fig. 7. The change in the slope of σ vs $^1/_T$ curves of $\text{Li}_{7\cdot1}\text{Ta}_{0\cdot9}\text{Zr}_{0\cdot1}0_6$ corresponds most likely again to the transition from extrinsic to intrinsic mechanism of conductivity. The activation energies are 0.50 eV and 0.60 eV in the lower and higher temperature regions respectively. A similar change in the slope of the σ vs $^1/_T$ plot of $\text{Li}_{7\cdot3}\text{Ta}_{0\cdot7}\text{Zr}_{0\cdot3}0_6$ is observed at a higher temperature, suggesting that the increased conductivity observed in the $\text{Li}_{7+x}\text{Ta}_{1-x}\text{Zr}_x0_6$ phases is due to the larger concentration of mobile Li^+ ions. Li_{7·4}Ta_{0.6}Zr_{0.4}0₆ is in the region of extrinsic conductivity for the entire range of temperature measured.

In Table 1, we summarize conductivity data for the best Li solid electrolytes. It may be concluded that solid solutions of lithium hexaoxometallates are also members of good lithium ion conducting solid electrolytes.

Fig. 7
The temperature dependence of conducitivity of Li_{7+X}Ta_{1-x}Zr_xO₆ phases

Table 1 Conductivity data for Li compounds.

Compounds	(Scm)-1	σ_{200}	Activation Energy (eV)	Temp. Range (°C)	Ref.
L1-β-Al ₂ 0 ₃	1.3 x 10 ⁻⁴	2.0×10^{-3}	0.19 0.36	-100 - 180 180 - 800	8
Li ₁₄ Zn(GeO ₄) ₄ (LISICON)		1.7 x 10 ⁻⁴	0.50	50 - 300	9
Li4B7012Cl0.68Br0.32	$5.4 \times 10^{-7*}$		0.53	50 - 230	1
Li3.75Si0.75P0.2504	4.8×10^{-7}			30 - 230	1
Li4.4Si0.6Al0.404	$2.8 \times 10^{-7*}$	7.6×10^{-4}	0.58	70 ~ 230	1
L14.6Al0.6S10.404		7.7×10^{-5}	0.68	85 - 230	1
Li7Ta0.7Nb0.306	7.3 x 10 ^{-8*}	4.3 x 10 ⁻⁴	0.63	80 - 230	this work
Li7Ta _{0.6} Bi _{0.4} 0 ₆	7.3 x 10 ^{-8*}	3.0 x 10 ⁻⁴	0.60	50 - 230	this work
Li _{7.4} Ta _{0.6} Zr _{0.4} 0 ₆	3.4 x 10 ^{-7*}	4.0 x 10 ⁻⁴	0.52	50 - 230	this work

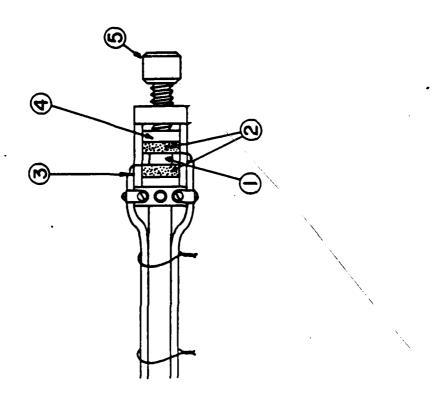
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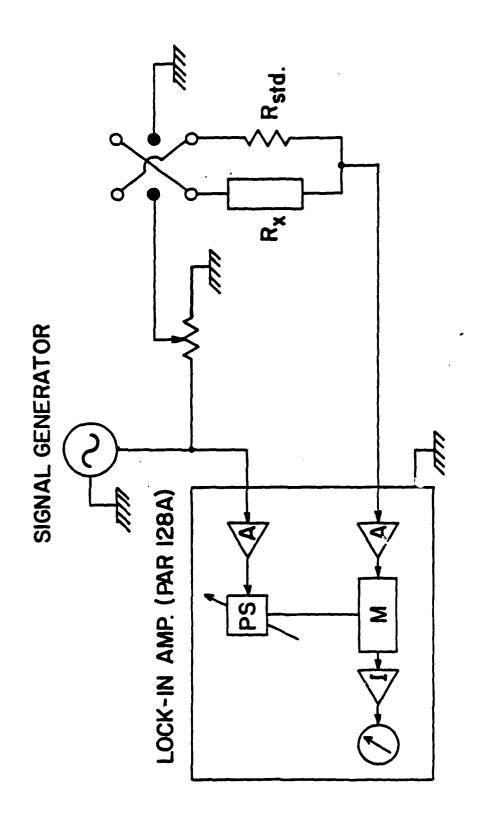
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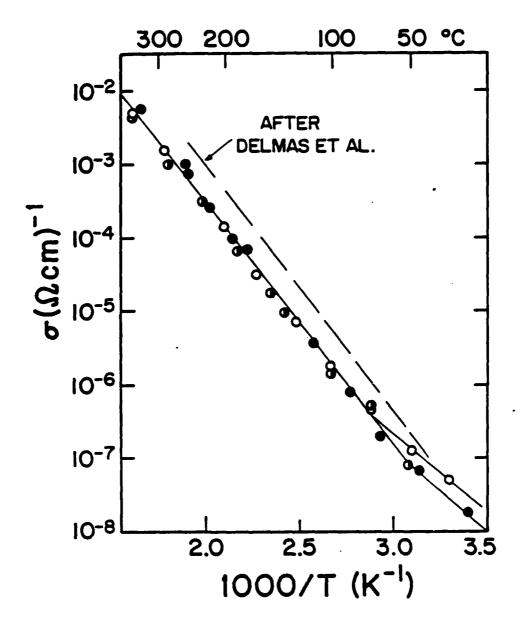
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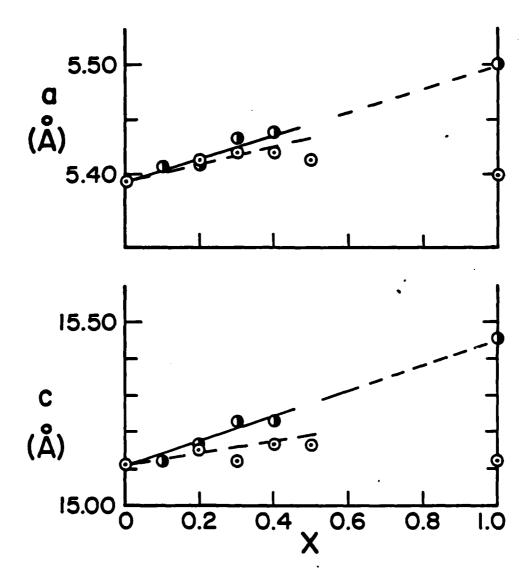
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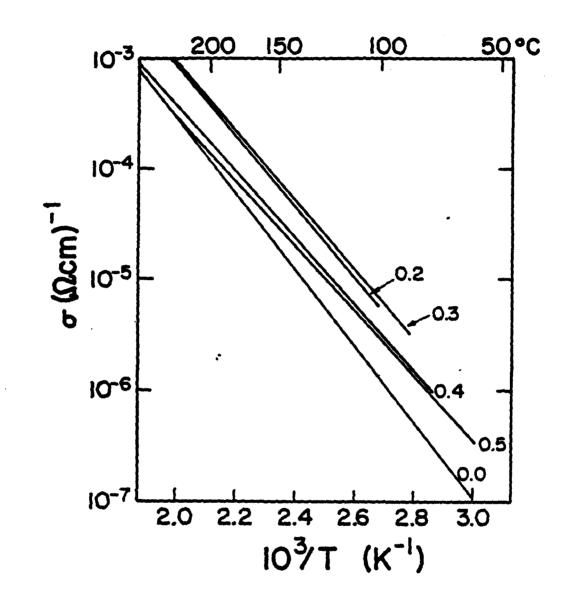


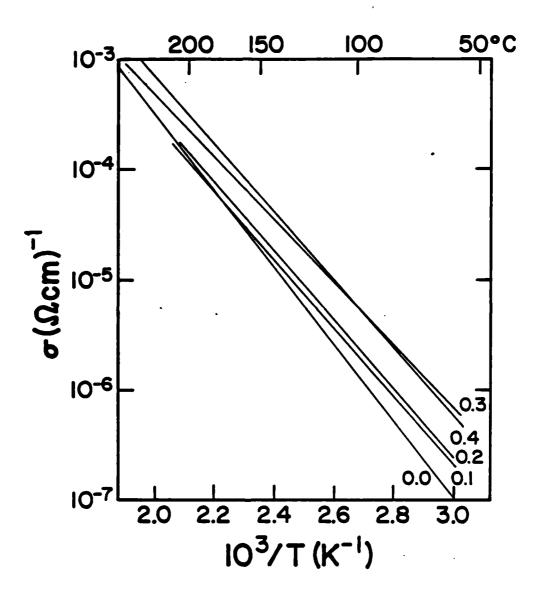
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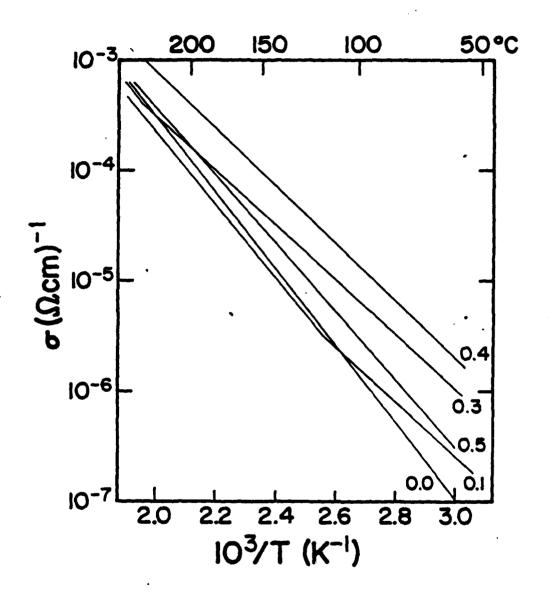












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